

Fuel-Composition Dependent Reactor Antineutrino Yield at RENO

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We report a fuel-dependent reactor electron antineutrino ($\bar{\nu}_e$) yield using six 2.8 GW_{th} reactors in the Hanbit nuclear power plant complex, Yonggwang, Korea. The analysis uses 850 666 $\bar{\nu}_e$ candidate events with a background fraction of 2.0% acquired through inverse beta decay (IBD) interactions in the near detector for 1807.9 live days from August 2011 to February 2018. Based on multiple fuel cycles, we observe a fuel ^{235}U dependent variation of measured IBD yields with a slope of $(1.51 \pm 0.23) \times 10^{-43} \text{ cm}^2/\text{fission}$ and measure a total average IBD yield of $(5.84 \pm 0.13) \times 10^{-43} \text{ cm}^2/\text{fission}$. The hypothesis of no fuel-dependent IBD yield is ruled out at 6.6σ . The observed IBD yield variation over ^{235}U isotope fraction does not show significant deviation from the Huber-Mueller (HM) prediction at 1.3σ . The measured fuel-dependent variation determines IBD yields of $(6.15 \pm 0.19) \times 10^{-43}$ and $(4.18 \pm 0.26) \times 10^{-43} \text{ cm}^2/\text{fission}$ for two dominant fuel isotopes ^{235}U and ^{239}Pu , respectively. The measured IBD yield per ^{235}U fission shows the largest deficit relative to the HM prediction. Reevaluation of the ^{235}U IBD yield per fission may mostly solve the reactor antineutrino anomaly (RAA) while ^{239}Pu is not completely ruled out as a possible contributor to the anomaly. We also report a 2.9σ correlation between the fractional change of the 5 MeV excess and the reactor fuel isotope fraction of ^{235}U .

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A definitive measurement of the smallest neutrino mixing angle θ_{13} has been a tremendous success in neutrino physics during the last decade [1,2]. The measurement has been achieved by comparing the observed $\bar{\nu}_e$ fluxes with detectors placed at two different distances from the reactors. As reactor $\bar{\nu}_e$ experiments suffer from large reactor related uncertainties of the expected $\bar{\nu}_e$ flux and energy spectrum [3–7], an identical detector configuration is essential to cancel out the systematic uncertainties. The reactor antineutrino anomaly (RAA), $\sim 6\%$ deficit of

measured $\bar{\nu}_e$ flux compared to the HM prediction is an intriguing mystery in current neutrino physics research and needs to be understood [4–6,8–11]. There have been numerous attempts to explain this anomaly by incorrect inputs to the fission β spectrum conversion, deficiencies in nuclear databases, underestimated uncertainties of reactor $\bar{\nu}_e$ model, and the existence of sterile neutrinos [3,12–19]. Moreover, all of the ongoing reactor $\bar{\nu}_e$ experiments have observed a 5 MeV excess in the IBD prompt spectrum with respect to the expected one [8,9,20,21]. This suggests that the reactor $\bar{\nu}_e$ model is not complete at all.

In commercial nuclear reactor power plants, almost all ($>99\%$) $\bar{\nu}_e$'s are produced through thousands of β -decay branches of fission fragments from ^{235}U , ^{239}Pu , ^{238}U , and ^{241}Pu . The $\bar{\nu}_e$ flux calculation is based on the inversion of spectra of the β -decay electrons of the thermal fissions which were measured in the 1980s at ILL [10,11].

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The reactor $\bar{\nu}_e$ models using these measurements as inputs have large uncertainties [5–7]. Therefore, reevaluation of the reactor $\bar{\nu}_e$ model and precise measurements of the neutrino flux and spectrum are essential to understand the RAA. Recently, Daya Bay Collaboration reported an observation of correlation between the reactor core fuel evolution and changes in the reactor $\bar{\nu}_e$ flux and energy spectrum [22]. The collaboration concluded that the ^{235}U fuel isotope may be the primary contributor to the RAA. In this Letter, we report an observation of a fuel-dependent variation of the reactor $\bar{\nu}_e$ flux using 1807.9 days of Reactor Experiment for Neutrino Oscillation (RENO) near detector data. We also present a hint of correlation between the 5 MeV excess and the reactor fuel isotope fraction of ^{235}U .

The Hanbit nuclear power plant complex consists of six reactor cores with total 16.8 GW_{th} in full operation mode. Two identical detectors are located at 294 (near detector) and 1383 m (far detector) from the reactor array center. The near (far) detector is under 120 (450) meters of water-equivalent rock overburden. The detectors with hydrocarbon liquid scintillator (LS) provide free protons as a target. Coincidence between a prompt positron signal and a delayed signal of gammas from neutron capture by Gadolinium (Gd) provides a distinctive IBD signature. Further details of the RENO detectors and $\bar{\nu}_e$ data analysis are found in Ref. [9].

The data used in this analysis are taken through IBD interactions in the near detector for 1807.9 live days from August 2011 to February 2018. Applying selection requirements yields 850666 IBD candidates with a background fraction of 2.0%. For the near detector data, we exclude a period of January to December 2013 because of detection inefficiency caused by an electrical noise coming from an uninterruptible power supply. We measure the reactor $\bar{\nu}_e$ flux as a function of an effective fission fraction $F_i(t)$ given by

$$F_i(t) = \frac{\sum_{r=1}^6 \frac{W_{th,r}(t) \bar{p}_r(t) f_{i,r}(t)}{L_r^2 \bar{E}_r(t)}}{\sum_{r=1}^6 \frac{W_{th,r}(t) \bar{p}_r(t)}{L_r^2 \bar{E}_r(t)}}, \quad (1)$$

where $f_{i,r}(t)$ is the fission fraction of the i th isotope in the r th reactor, $W_{th,r}(t)$ is the r th reactor thermal power, $\bar{p}_r(t)$ is the mean survival probability of $\bar{\nu}_e$ from the r th reactor, and L_r is the distance between the near detector and the r th reactor. The average $\bar{\nu}_e$ energy produced following fission to be converted into heat is given by $\bar{E}_r(t) = \sum_{i=1}^4 f_{i,r}(t) \langle E_i \rangle$, where $\langle E_i \rangle$ is an average energy released per fission where $(\langle E_{235} \rangle, \langle E_{238} \rangle, \langle E_{239} \rangle, \langle E_{241} \rangle) = (202.36, 205.99, 211.12, 214.26)$ energy per fission/MeV [23]. The upper panel of Fig. 1 shows time variation of the effective fission fraction of ^{235}U viewed by the near detector. The effective fission fraction is obtained from the daily thermal power and fission fraction data of each reactor core, provided by the Hanbit nuclear power plant. A total average IBD yield (\bar{y}_f) is

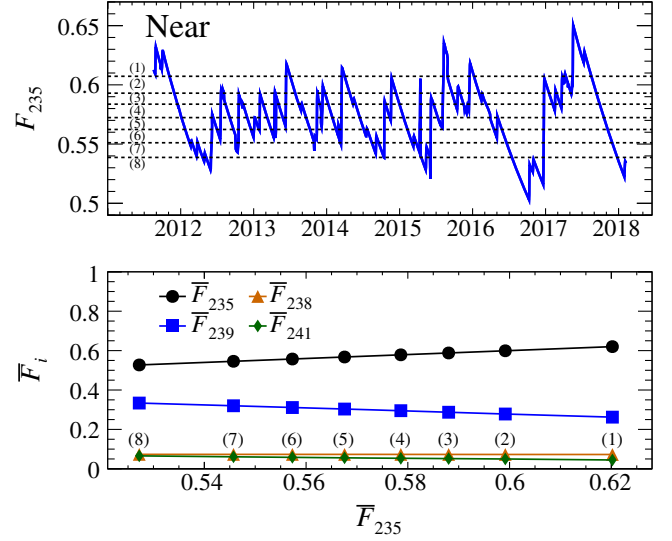


FIG. 1. Top: Effective ^{235}U daily fission fraction (F_{235}) in the near detector according to Eq. (1). The daily F_{235} is obtained from the reactor information provided by the Hanbit nuclear power plant. Bottom: Relative fission fractions for the primary fuel isotopes of ^{235}U , ^{239}Pu , ^{238}U , and ^{241}Pu . The numbers in the parentheses represent eight data groups with different fission fractions [24].

measured to be $\bar{y}_f = (5.84 \pm 0.13) \times 10^{-43} \text{ cm}^2/\text{fission}$ with average effective fission fractions F_{235} , F_{238} , F_{239} , and F_{241} of 0.573, 0.073, 0.299, and 0.055, respectively.

For examining fuel-dependent variation of reactor $\bar{\nu}_e$ yield, eight groups of equal data size are sampled according to the eight different values of the ^{235}U fission fraction. A time-averaged effective fission fraction ($\bar{F}_{i,j}$) of the i th isotope in the j th data group is calculated as

$$\bar{F}_{i,j} = \frac{\int dt \sum_{r=1}^6 \frac{W_{th,r}(t) \bar{p}_r(t) f_{i,r}(t)}{L_r^2 \bar{E}_r(t)}}{\int dt \sum_{r=1}^6 \frac{W_{th,r}(t) \bar{p}_r(t)}{L_r^2 \bar{E}_r(t)}}. \quad (2)$$

The time-averaged effective fission fractions of the four isotopes in each group are shown as a function of time-averaged fission fraction of ^{235}U (\bar{F}_{235}) in the lower panel of Fig. 1. An average IBD yield per fission of the j th data group ($\bar{y}_{f,j}$) is given by

$$\bar{y}_{f,j} = \sum_{i=1}^4 \bar{F}_{i,j} y_i, \quad (3)$$

where an integrated IBD yield per fission (y_i) is calculated as $y_i = \int \sigma(E_\nu) \phi_i(E_\nu) dE_\nu$, $\sigma(E_\nu)$ is the IBD reaction cross section, and $\phi_i(E_\nu)$ is the reactor $\bar{\nu}_e$ spectrum generated by each reactor's fission isotope, $(y_{235}, y_{239}, y_{238}, y_{241}) = (6.70 \pm 0.14, 4.38 \pm 0.11, 10.07 \pm 0.82, 6.07 \pm 0.13) \times 10^{-43} \text{ cm}^2/\text{fission}$ [7]. We use the IBD cross section in Refs. [7,25] and a neutron lifetime of 880.2 s

in the calculation [26]. The IBD yield y_i of a fissile isotope is sensitive to its reactor $\bar{\nu}_e$ spectrum because the IBD cross section increases with the $\bar{\nu}_e$ energy. A model-independent IBD yield of $\bar{y}_{f,j}$ is determined by counting the number of events in each data group using the following relationship.

$$N_j = \bar{y}_{f,j} \sum_{r=1}^6 \frac{N_p}{4\pi L_r^2} \int \left(\frac{W_{th,r}(t) \bar{P}_r(t)}{\sum_i f_{i,r}(t) \langle E_i \rangle} \right) \epsilon_d(t) dt, \quad (4)$$

where N_j is the number of IBD events in the j th group, N_p is the number of target protons, $\bar{P}_r(t)$ is the mean survival probability, and $\epsilon_d(t)$ is the detection efficiency including the signal loss due to timing veto requirements. The average IBD yield of $\bar{y}_{f,j}$ for each data group is determined by the observed N_j . No fission-fraction dependent IBD yield expects a flat distribution of \bar{y}_f as a function of \bar{F}_{235} . There are several updates in this analysis from the previous publication [27]. They are use of an IBD cross section in Ref. [25], an updated detection efficiency including the neutron spill-out effect, and an improved thermal energy release per fission in Ref. [23]. A detailed description of the updates will be reported in an upcoming publication. Figure 2 shows a measured distribution of \bar{y}_f as a function of \bar{F}_{235} or \bar{F}_{239} for the eight data groups. We observe a clear correlation between \bar{y}_f and \bar{F}_{235} , indicating dependence of the IBD yield per fission on the isotope fraction of ^{235}U . A linear function is used for a fit to the eight data points with $\chi^2/\text{NDF} = 4.60/6$ at the best fit. The horizontal line represents an expected distribution for no fuel-dependent IBD yield. This result rules out no fuel-dependent variation of the IBD yield per fission at 6.6σ confidence level,

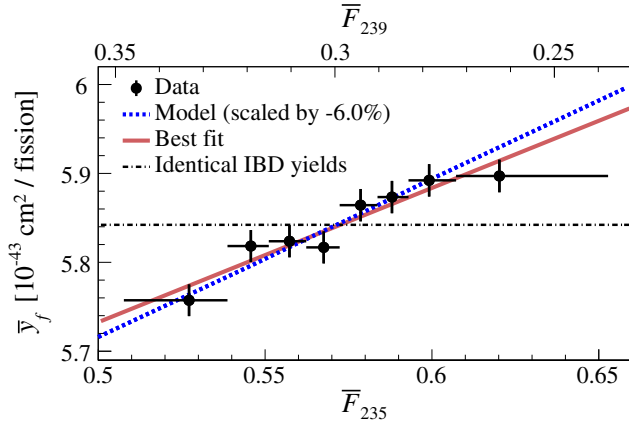


FIG. 2. IBD yield per fission \bar{y}_f as a function of the ^{235}U effective fission fraction. The measured values (black dots) are compared to the scaled HM prediction (blue dotted line) and the best fit of the data (red solid line). The value of \bar{F}_{235} for each data point is calculated as an average of fission fractions weighted by thermal power and a distance between reactor and detector. The error of \bar{F}_{235} indicates the variation of ^{235}U fission fraction. Errors of \bar{y}_f are statistical uncertainties only [24].

corresponding to the p value of 3.4×10^{-11} . It indicates that the variation of the \bar{y}_f as a function of \bar{F}_{235} comes from unequal IBD yields among different isotope fissions. The measured yield variation is fitted with the HM prediction to obtain the best fit at a scaling of -6.0% with $\chi^2/\text{NDF} = 6.25/7$. Thus the observed IBD yield variation over ^{235}U fission fraction is not inconsistent with the HM prediction at 1.3σ . The measured IBD yield variation is also fit with the prediction from the *ab initio* calculation in Ref. [28]. A best-fit of $\chi^2/\text{NDF} = 4.79/7$ is found at a scaling of -5.1% to make a better agreement with the data in the slope.

For determination of y_{235} and y_{239} simultaneously, a χ^2 with pull parameter terms of systematic uncertainties is constructed using the observed IBD yield per fission and minimized by varying the free parameters of y_{235} and y_{239} , and pull parameters. The subdominant isotopes of ^{238}U and ^{241}Pu are constrained in the fitter within uncertainties of 10% [4] and 5% [29], respectively. The uncertainties of thermal power, fission fraction, energy per fission, and detection efficiency are considered to be fully correlated among the eight data groups in the different fission fraction bins. Each correlated uncertainty is taken into account through a pull parameter in the χ^2 calculation. The χ^2 is given by

$$\chi^2 = \sum_{j=1}^8 \left(\frac{\bar{y}_{\text{obs},j} - \bar{y}_{\text{exp},j}}{\sigma_{\text{obs},j}} \right)^2 + \left(\frac{\xi_{238}}{\sigma_{238}} \right)^2 + \left(\frac{\xi_{241}}{\sigma_{241}} \right)^2 + \left(\frac{\xi_{\text{th}}}{\sigma_{\text{th}}} \right)^2 + \left(\frac{\xi_{\text{f}}}{\sigma_{\text{f}}} \right)^2 + \left(\frac{\xi_{\text{en}}}{\sigma_{\text{en}}} \right)^2 + \left(\frac{\xi_{\text{det}}}{\sigma_{\text{det}}} \right)^2, \quad \text{where}$$

$$\bar{y}_{\text{exp},j} = [\bar{F}_{235}^j y_{235} + \bar{F}_{239}^j y_{239} + \bar{F}_{238}^j y_{238} (1 + \xi_{238}) + \bar{F}_{241}^j y_{241} (1 + \xi_{241})] (1 + \xi_{\text{th}} + \xi_{\text{f}} + \xi_{\text{en}} + \xi_{\text{det}}), \quad (5)$$

where $\bar{y}_{\text{obs},j}$ is the observed IBD yield per fission averaged over the four isotopes in the j th data group, $\sigma_{\text{obs},j}$ is the statistical uncertainty of $\bar{y}_{\text{obs},j}$, $\bar{y}_{\text{exp},j}$ is the expected IBD yield per fission averaged over the four isotopes, \bar{F}_i^j is the time-averaged effective fission fraction of the i th isotope for the j th data group, σ_{238} and σ_{241} are the uncertainties of y_{238} (10%) and y_{241} (5%), respectively, σ_{th} , σ_{f} , σ_{en} and σ_{det} are the uncertainties of thermal power (0.5%), fission fraction (0.7%), energy per fission (0.2%) and detection efficiency (1.93%), respectively. The correlated uncertainties among the eight data groups are considered by changing pull parameters in the $\bar{y}_{\text{exp},j}$ within their uncertainties. Each pull parameter is common among the eight data groups to treat its fully correlated uncertainty. ξ_{238} and ξ_{241} are the pull parameters of y_{238} and y_{241} , respectively, and ξ_{th} , ξ_{f} , ξ_{en} , and ξ_{det} are the pull parameters for thermal

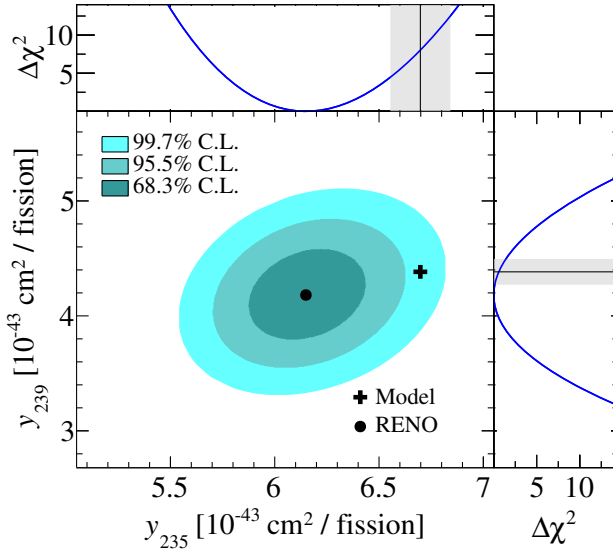


FIG. 3. Combined measurement of y_{235} and y_{239} . The shaded contours are allowed regions and the dot is the best fit. The cross shows the prediction of the HM. The top and right side panels show one dimensional $\Delta\chi^2$ profile distributions for y_{235} and y_{239} while the gray shaded bands represent the model predictions.

power, fission fraction, energy per fission, and detection efficiency, respectively.

The best-fit results are $y_{235} = (6.15 \pm 0.19) \times 10^{-43}$ and $y_{239} = (4.18 \pm 0.26) \times 10^{-43}$ cm²/fission. Figure 3 shows the combined measurement of y_{235} and y_{239} . The measured IBD yield per ^{235}U fission is smaller than the HM prediction at 2.8σ while the measured yield per ^{239}Pu fission is smaller than the prediction only at 0.8σ . This suggests that the RAA

can be largely understood by incorrect estimation of the ^{235}U IBD yield.

Following the analysis in Ref. [30] we also perform the combined measurements for all combinations of the four isotopes, total six pairs. The χ^2 of Eq. (5) is used with an added constraint term of $(\xi_i/\sigma_i)^2$, where σ_i are uncertainties of y_{235} (5%), y_{238} (10%), y_{239} (5%), and y_{241} (5%) [30]. Figure 4 shows the allowed regions of each pair of IBD yields per fission. The dot is the best fit of each pair of IBD yields while the crossing lines represent the HM predicted yields. The shaded contours are 68.3, 95.5, and 99.7 % C.L. allowed regions for each pair of IBD yields. In the fitting results of the six pairs of isotopes, we observe that y_{235} is smaller than the prediction at $\sim 2.5\sigma$ while the IBD yields per fission of the rest isotopes are consistent with the prediction within 1σ .

The deficit of y_{235} relative to the HM prediction could be interpreted by an indication of incorrectly evaluated IBD yield of ^{235}U fission that may be a major source of the RAA [22,31]. As the ^{235}U and ^{239}Pu fission fractions are correlated, we perform pseudoexperiments to test this possibility. Pseudodata with IBD yields per fission, $\bar{y}_{f,j}$, are produced for various ratios of y_{235} and y_{239} . For each input of y_{235} and y_{239} , 1000 pseudodata are produced within statistical errors. In addition, pseudoexperiments with \bar{y}_f scaled down by 6.0% from the model prediction are generated by reducing y_{235} only. A fit finds a value of y_{235} less than the measured value with 3.4σ deviation from the model prediction. This does not reproduce the measured y_{235} deviation of 2.8σ by reducing y_{235} only, while a pseudoexperiment of y_{235} down by 8.2% and y_{239} down by 4.6% reproduce the observed

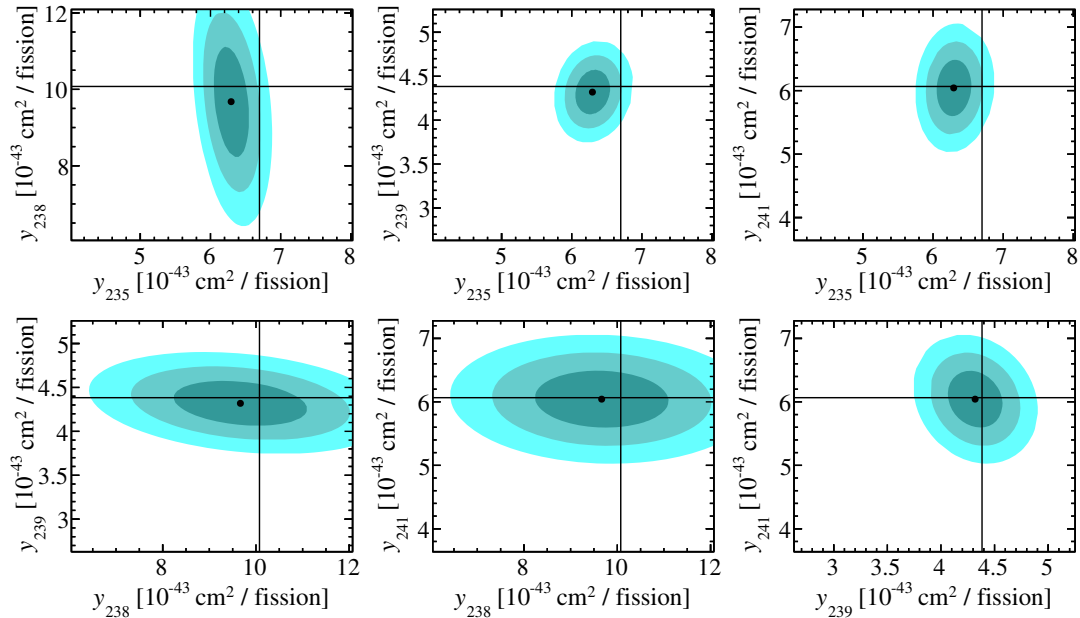


FIG. 4. Allowed regions of the IBD yield per fission for the six pairs of fission isotopes. The dots indicate the best-fit IBD yields and the cross lines represent the model prediction. The three contours are allowed regions of 68.3, 95.5, and 99.7 % C.L.

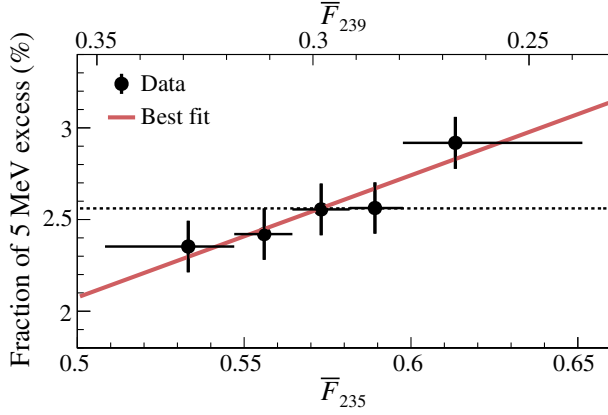


FIG. 5. Fraction of the 5 MeV excess as a function of \bar{F}_{235} . The red line is the best fit to the data and the dotted line represents no correlation of 5 MeV excess fraction with \bar{F}_{235} [24].

deviations of 2.8σ and 0.8σ , respectively. Thus, we conclude that RENO data do not rule out ^{239}Pu as a contributor to the RAA and the anomaly can be mostly explained by reevaluation of the ^{235}U IBD yield.

The RENO collaboration has reported an excess of the observed IBD prompt spectrum at 5 MeV [8,9], also observed by the other ongoing reactor $\bar{\nu}_e$ experiments as well [20,21]. The 5 MeV excess is observed to be proportional to the reactor thermal power [9]. Several explanations and suggestions are proposed to understand the origin of the 5 MeV excess [12–14,32,33]. There is a suggestion that a particular isotope could be the source of the excess [13], while an analysis disfavors the ^{239}Pu and ^{241}Pu isotopes as a single source for the 5 MeV excess [34]. However, there is no clear understanding of the origin of the 5 MeV excess yet.

A possible fuel dependence of the 5 MeV excess is examined by the IBD yield per fission for the events only in the 5 MeV region of $3.8 < E_p < 7$ MeV. We have not seen any significant deviation of the IBD yield slope with respect to the HM prediction between the 5 MeV region and the entire energy range of 1.2 to 8 MeV, consistent with the results from Daya Bay [22]. For a fuel dependence of the 5 MeV excess only, an event rate with the HM expected energy spectrum in the 5 MeV region is obtained by a fit to the data in the energy ranges of 1.2 to 3.8 MeV and 7.0 to 8.0 MeV and subtracted from the total 5 MeV rate. Five groups of equal data size are sampled according to five different values of \bar{F}_{235} . A fraction of the 5 MeV excess is calculated as a ratio of the 5 MeV excess rate to the total IBD rate in the entire energy range. Figure 5 shows the distribution of 5 MeV excess fractions as a function of \bar{F}_{235} . The best fit for the data with a first-order polynomial function shows a correlation between the 5 MeV excess fraction and \bar{F}_{235} . The data are also fitted with a zeroth-order polynomial function with an average excess fraction of $(2.56 \pm 0.06)\%$. The hypothesis of no correlation between the 5 MeV excess fraction and \bar{F}_{235} is disfavored

at 2.9σ where the χ^2/NDF is 1.17/3 for the best fit and 9.58/4 for no-correlation hypothesis. While the current result shows an indicative correlation of the 5 MeV excess fraction with \bar{F}_{235} and an anticorrelation with the rest isotope fractions, further accumulated data may reveal the source of the 5 MeV excess. We repeat extraction of the 5 MeV excess by subtracting the HM prediction estimated from reactor thermal powers and fuel isotope fractions. The significance of correlation between the 5 MeV excess fraction and \bar{F}_{235} becomes 1.3σ . The data-driven subtraction described earlier is free from the uncertain HM flux normalization.

In summary, we report a fuel-dependent IBD yield using 1807.9 days of RENO near detector data. We measure IBD yields per fission of $(6.15 \pm 0.19) \times 10^{-43}$ and $(4.18 \pm 0.26) \times 10^{-43} \text{ cm}^2/\text{fission}$ for the dominant fission isotopes of ^{235}U and ^{239}Pu , respectively. A change in the IBD yield with respect to the effective ^{235}U fission fraction is observed at 6.6σ . The measured IBD yield per fission of $(5.84 \pm 0.13) \times 10^{-43} \text{ cm}^2/\text{fission}$ is 6.0 % smaller than the HM prediction and confirms the RAA. The measured IBD yield per ^{235}U fission is smaller than the HM prediction at 2.8σ . This suggests that the RAA can be largely understood by incorrect estimation of the ^{235}U IBD yield. We obtain the first hint (2.9σ) for a correlation between the 5 MeV excess fraction and the ^{235}U fission fraction.

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